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MEMORANDUM REPORT NO. 2509

CALCULATIONS OF FRAGMENT VELOCITIES FROM NATURALLY FRAGMENTING MUNITIONS

R. R. Karpp W. W. Predebon

July 1975

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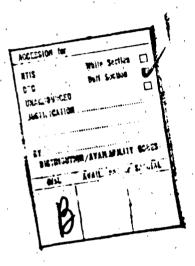
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I. INTRODUCTION

Accurate modeling as well as predicting the performance of naturally fragmenting munitions is particularly important to the Army munitions community. For parametric studies, calculational techniques which are both quick and easy to use are of practical interest. One well known technique employs the Gurney¹ and Taylor² formulae. In Section II the advantages and disadvantages of this technique are demonstrated.

In Section III the importance of gas leakage after casing breakup is discussed with the use of the LLL standard cylinder test³ data and experimental observations of a variety of test warhead firings.

In Section IV the use of a time-dependent, two dimensional, finite-difference code, HEMP, to predict fragment velocities is discussed in detail. Problems associated with determining a definite final fragment velocity are mentioned. In an effort to model casing acceleration realistically, a gas leakage model is presented as an integral part of the code calculations. Lastly, the code calculations are compared with a number of experimental test results.

II. APPLICABILITY OF THE GURNEY AND TAYLOR FORMULAE

Practically speaking, it is desirable to have a method available for estimating fragment velocities from naturally fragmenting munitions which is both quick and easy to use for parametric studies. One such method for estimating fragment velocity is the Gurney formula (Figure 1) which gives the fragment speed as a function of the charge to metal ratio (mass of explosive/mass of casing) and an empirically determined constant. The direction that the fragment assumes can then be determined by the Taylor formula (Figure 1) using this calculated speed. For an axisymmetric explosive-metal system with variable diameter, the charge to metal ratio (C/M) is computed locally as a function of axial position. Thus, the speed computed from the Gurney formula varies with axial position. In the Taylor formula for the projection angle, both the fragment speed and the rate at which the detonation wave traverses the casing varies with axial position.

^{1.} Gurney, R.W., "The Initial Velocities of Fragments From Bombs, Shells and Grenades". BRL Report 405. September 1943. AD #36218

^{2.} Birkhoff, G., MacDougall, P., Pugh, E.M., and Taylor, G.I., Sir., "Explosives With Lined Cavities", J. Appl. Phys., Vol. 19, p. 563, June 1948.

^{3.} Lee, E.L., Hornig, H.C., and Kury, J.W., "Adiabatic Expansion of High Explosive Detonation Products", UCRL-50122, May 2, 1968.

Numerous comparisons between calculated fragment velocities using the Gurney and Taylor formulae and experimental test results have been made by us and other researchers. For a practical range of charge to metal ratios (.1 < C/M < 2), it is our observation that speeds and directions thus calculated agree reasonably well with experimental data for configurations when the flow is essentially radial. Figure 2 illustrates the results of a comparison of the above technique with experimental data for a typical H.E. projectile. The agreement with this test data is quite good. A second example of the agreement between the results of this technique and experimental data is shown in Figure 3. Here a steel cylinder with length to diameter ratio of two filled with Composition B explosive and initiated on axis at one end is shown. Here the agreement is not good. This figure illustrates the degree to which end effects may influence the fragment speed distribution and projection angle distribution. One can see that a method which can predict both the fragment speed and direction distributions correctly is desirable. That is, a method which will model general time-dependent, two-dimensional phenomena, such as end effects, with some degree of confidence would be a useful tool. The method proposed and utilized in this report is a modified version of the time-dependent, two-dimensional, finite-difference computer code, HEMP.

III. EFFECT OF GAS LEAKAGE AFTER CASING BREAKUP

The standard cylinder test³ developed by the Lawrence Livermore Laboratory, Livermore, California, provides much data concerning the expansion of explosive-filled cylinders. Figure 4 schematically illustrates the test procedure. The radial expansion of the copper cylinder is measured with a streak camera at a point two-thirds the length of the cylinder from the initiated end. The rate of expansion of the cylinder is plotted in Figure 5 as a function of radial displacement and expansion ratio for three different explosive fills, Octol, Composition B and TNT. The percentages indicated on the curves signify the percent of the final expansion rate. During these measurements, neither significant leakage of explosive products nor influence of end effects occurs. The general shape of the curves are of interest since they indicate that most of the acceleration has occurred by an expansion ratio of R/R₀ = 2. Even at R/R₀ = 1.75 at least 92% of the final velocity has already been obtained. These curves are for copper casings; however, Figure 6 (taken from Reference 3) indicates that the shape of the curve is the same for stainless steel and mild steel; the latter is important for naturally fragmenting, steel casing munitions.

Gas leakage has been experimentally observed at an expansion ratio of about $R/R_0=1.75$ for many steels.^{4,5} Figure 7 lists such observed values of R/R_0 for steel, copper, and aluminum casings. Even though about 92% or more of the final casing velocity is acquired before casing breakup for natural fragmentation; we will show in Section IV that it is desirable to include the effects of gas leakage in time dependent calculations to obtain accurate predictions of fragment velocities particularly in regions influenced by end rarefaction waves.

- IV. USE OF TIME-DEPENDENT, TWO-DIMENSIONAL, FINITE-DIFFERENCE CODES
 - 1. Equation of State of the Explosive

The Jones, Wilkins, Lee (JWL) equation of state recommended in Reference 3 and given in Figure 8 is used in these calculations. The finite-difference code used in the present work is the HEMP code⁶. However, equally good results could be obtained with other valid finite-difference codes⁷.

The JWL equations of state are developed such that the radial expansion rate obtained from the LLL standard cylinder test can be duplicated with a HEMP calculation. Figure 8 indicates the accuracy with which the finite-difference calculations reproduce the data. Figure 9 shows the finite-difference mesh which was used to compare with the LLL standard cylinder test data. In general, good agreement between HEMP code calculations, experiments, and exact analytical solutions has been demonstrated.⁸

2. Problems Associated With Determination of a Final Fragment Velocity

Figure 10 shows the finite-difference mesh at various times for a HEMP calculation of an open ended steel cylinder with an Octol

^{4.} Taylor, G.1., "The Fragmentation of Tubular Bombs", <u>Scientific</u> Papers of G.1. Taylor, Vol. III, Cambridge University Press, London, 1963.

^{5.} Hoggatt, C. and Recht, R., "Fracture Behavior of Tubular Bombs", J. Appl. Phys., Vol. 39, pp. 1856 - 1862, February 1968.

^{6.} Wilkins, M.L., "Calculation of Elastic-Plastic Flow", Methods of Computational Physics, Vol. 3, edited by Alder, B., Fernbach, S. and Rotenburg, M., Academic Press, New York, 1964.

^{7.} Sedgwick, R.T., A.J. Good, and L.J. Hageman, "Theoretical Investigation of High Explosive Fragmentation Munitions", Final Report, Contract DAAD 05-71-C-0092, Systems, Science, and Software Report 3SR-867, Nov. 24, 1971.

^{8.} Karpp, R.R., "Accuracy of HEMP Code Solutions", BRL Memo Report 2268, January 1973. AD #757153

explosive fill. The steel casing can be modeled in several ways. way is to treat the casing as a fluid (neglect the strength of the steel). This method produces the velocity* - time curves shown in Figure 11 (solid curve). The early portion of the acceleration curve is essentially correct. However, since it is known that a steel cylinder will fragment at an expansion ratio of about 1.75, the latter portion of the acceleration curve will tend to over estimate the final fragment velocity. This effect is a result of the code requiring the casing to be continuous, and, therefore, the internal pressure acts on a continuously increasing area and produces an unrealistically high fragment velocity. Referring to Figure 11, this effect is more pronounced at position 1, where end effects are important, than at position 2. Therefore, the calculational results with the fluid model will tend to overestimate the final fragment velocity. Since the velocity continually increases, one must use an artificial method to determine the correct final fragment velocity distribution. A second method of modeling the casing includes the material strength of the casing. This can be done by modeling the metal casing as an elastic-plastic material. A curve showing the velocity-time plot when the casing is treated as an elasticperfectly plastic material is given in Figure 11 (dot-dash curve). This model does produce a more definite valu, for the final fragment velocity. However, this maximum value for fragmert 'elocity still occurs at an' expansion ratio considerably beyond the known expansion ratio at fragmentation. Besides the fluid model objection which also applies to this model, it does not seem realistic to include strength in the calculations past the point where the casing fragments. A calculational procedure which does not rely on an artificial method to determine the final fragment velocity but includes the ap ropriate physics of fragmentation and subsequent gas leakage is discussed in the next section.

3. Final Fragment Velocities Using a Gas Leakage Model

Several features of the warhead expansion and fracture process are inadequately modeled by axisymmetric finite-difference computer codes when casing breakup occurs. First, the increase in circumference ceases when casing fragmentation occurs. The standard from of these codes causes the casing to remain continuous. This treatment tends to yield a higher radial acceleration and thus higher final fragment velocity since the pressure acts on an area which is larger than the true area. However, this effect can be corrected by calculating the true force acting on an element of the casing. After casing fragmentation (R>Rb) the true force is calculated by reducing the pressure acting on the interface of the casing by the factor $R_{\rm b}/R$ (Refer to Figure 12).

^{*} Velocity as used in conjunction with any of the Figures given is the magnitude of the velocity. When the direction of the velocity is discussed or plotted it is referred to as projection angle.

Secondly, after casing fragmentation, i.e. $R > R_{\rm b}$, a void area $A_{\rm L}$ opens up and gas leakage occurs. This effect is modeled by calculating the rate of efflux from an ideal nozzle passing its maximum rate of flow. For example, for an ideal gas the mass rate of flow is given by

$$\dot{m} = A_L \sqrt{\gamma P \rho} \left(\frac{2}{\gamma + 1} \right)^{\frac{\gamma + 1}{2(\gamma - 1)}}$$

Where P is pressure, ρ is density and γ is the ratio of specific heats. This rate of efflux is calculated in the model for every cycle and the mass leaked for a single cycle is $\Delta m = m \Delta t$. The amount of mass (Δm) is then extracted from the zone in contact with the casing interface. In this manner, the model accounts for gas leakage around the fragments after casing breakup. Thirdly, to simulate a loss in circumferential stress upon casing breakup, the yield strength of the casing material is set equal to zero when $R > R_b$.

A comparison of the three casing models discussed, i.e. fluid, elastic-plastic and elastic-plastic with gas leakage, is shown in Figure 11. Here the velocity-time curves are given for two positions along the casing, namely, close to the initiation end and at the maximum velocity position. Figure 11 clearly shows that a more definite and lower final fragment velocity is obtained by modeling the casing material as an elastic-plastic material and accounting for gas leakage.

4. Comparisons of Code Calculations with Experimental Results

In this section we present comparisons of the code calculations, utilizing the gas leakage model, with a number of experimental test results to demonstrate the reliability of this calculational technique. In Figures 13 through 18 the HEMP code predictions of fragment speed and projection angle are compared with experimental test results. The HEMP code calculations of fragment speed and projection angle are plotted as continuous curves in all the figures. The experimental data are plotted according to round numbers in each case. For information concerning the experimental set up, data collection and test firings, the reader is referred to Reference 9 for details. For each of the six cases represented in Figures 13 through 18 the length to diameter ratio was two, and they were end initiated as shown in Figure 13. The explosive fills were Octol, Composition B, and TNT, and the casing was 1020 steel or HF-1 steel. The charge to metal ratio

^{9.} Karpp, R.R., Kronman, S., Dietrich, A.M. and Vitali, R., "Influence of Explosive Parameters on Fragmentation", BRL Memorandum Report 2336, October 1973. AD #917248L

was 0.8 (thin wall) for the cases shown in Figures 13-15 and 0.4 (thick wall) in Figures 16 - 18. No discussion will be given here concerning the experimental data as it is covered in detail in Reference 9. However, we do wish to emphasize the good agreement between experimental results and the calculations for a variety of explosive fills and C/M ratios.

Lastly, in Figures 19 and 20 HEMP calculations are again compared with test data. However, in this case the test data are from "arena tests" of an H.E. projectile. ¹⁰ In Figure 19 the fragment velocity is plotted against angular position in the arena test set up. In Figure 20 the percent recovered mass in each angular recovery position is plotted. Again, very good agreement is demonstrated.

V. SUMMARY AND RECOMMENDATIONS

In the first part of the report it was shown that fragment velocity predictions based on the Gurney and Taylor formulae may be adequate for cases when the flow is one-dimensional (radial), but inadequate when the flow becomes two-dimensional. It was then shown from the LLL standard cylinder tests that in the natural fragmentation case most of the casing velocity is obtained prior to fragmentation. Next the use of a time-dependent, two-dimensional, finite-difference code, HEMP, to mode! naturally fragmenting munitions was discussed. Arguments were given to show that the fluid model (without strength) for the casing material will tend to over-estimate the final fragment velocity. The casing material modeled as an elastic-plastic material tends to produce a more definite final fragment velocity, but at unrealistically high cylinder expansions. In an effort to remove some of these objections and at the same time model the casing motion more realistically, a gasleakage model was proposed in conjunction with the code calculations. The gas leakage model was designed to simulate explosive gas leakage around fragments after casing breakup. Comparisons are given between the code calculations with the gas leakage model and experimental data, and very good agreement was demonstrated. However, we wish to emphasize that the gas leakage model presented is still in its early development. Further modifications and improvements are planned. Based on the agreement shown between the code calculations and the experimental test cases given, it is our judgement that this calculational tool can be used confidently in a variety of naturally fragmenting munitions when experimental data is not available.

Our future efforts will include the development of a more general criterion to initiate the gas leakage based on the state of stress

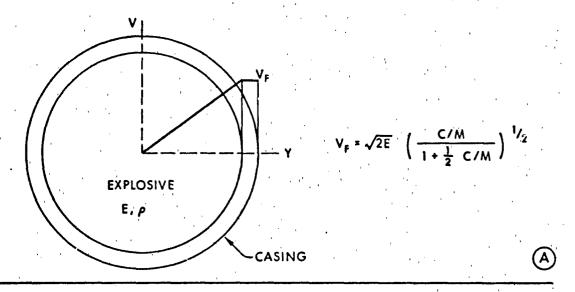
^{10.} U.S. Army TECOM Development and Proof Services, Firing Record No. B11982, Aberdeen Froving Ground, Maryland, August 1953 to June 1954.

in the casing material, rather than on the expansion ratio. Lastly, some effort has already been expended and much more is planned to investigate the use of the HEMP calculations to model fragment velocities from preformed land controlled fragmentation warheads using the proposed gas leakage model.

^{71.} Karpp, R.R., and Predebon, W.W., "Calculation of Fragment Velocities from Fragmentation Munitions" First International Symposium on Ballistics, Orlando, Flz., 13-15 November 1974, Section IV, pp. 145-176, American Defense Preparedness Association, Washington, D.C.

VI. APPENDIX A: FIGURES

GURNEY'S VELOCITY FORMULA



TAYLOR'S PROJECTION ANGLE FORMULA

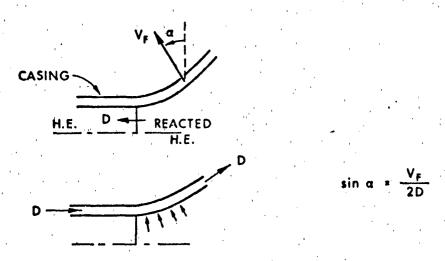


Figure 1. Gurney Velocity Formula (Cylinder) and Taylor's Projection Angle Formula

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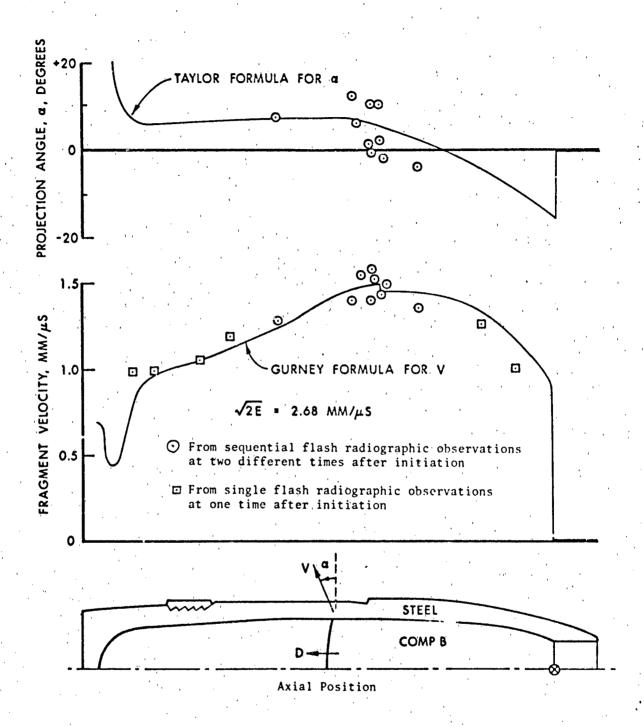
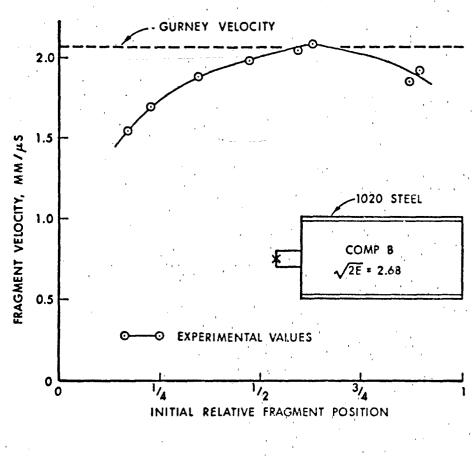


Figure 2. Comparison of Gurney and laylor Formulae Predictions With Experimental Data, Case 1



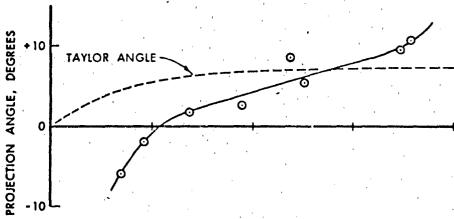


Figure 3. Comparison of Gurney and Taylor Formulae Predictions With Experimental Data, Case 2

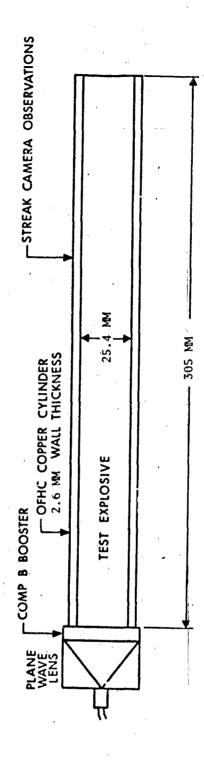


Figure 4. LLL Standard Cylinder Test



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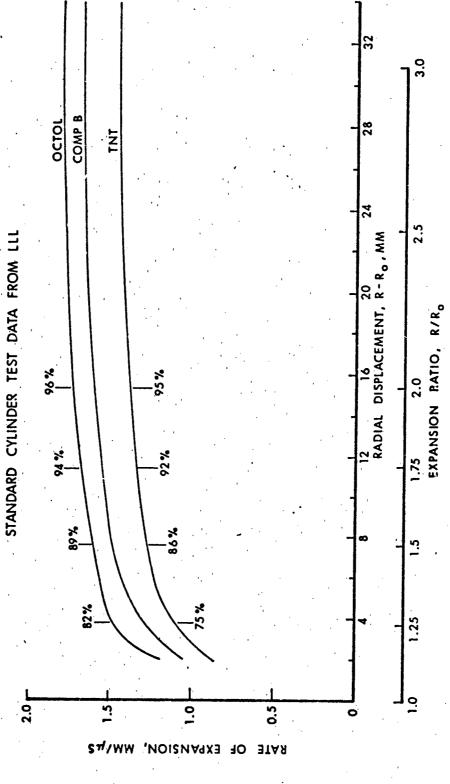


Figure 5. Comparison of the Rate of Expansion Versus Expansion Ratio for Various Explosive Fills (Ref. 3)

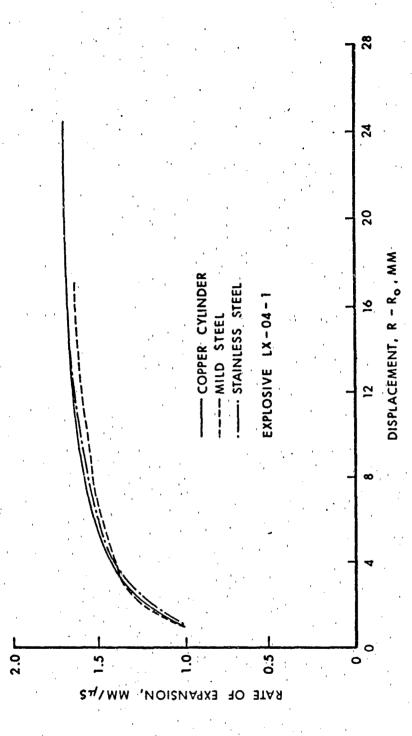
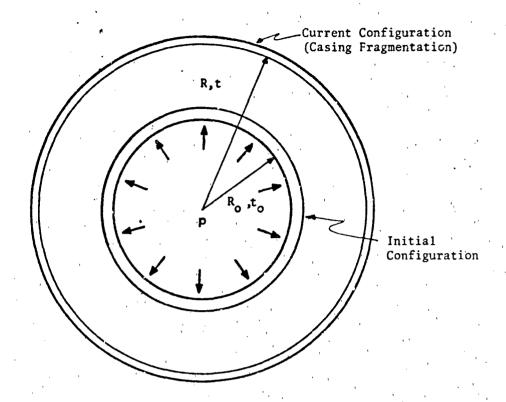


Figure 6. Comparison of Rate of Expansion Versus Displacement for Various Casing Materials (Ref. 3)



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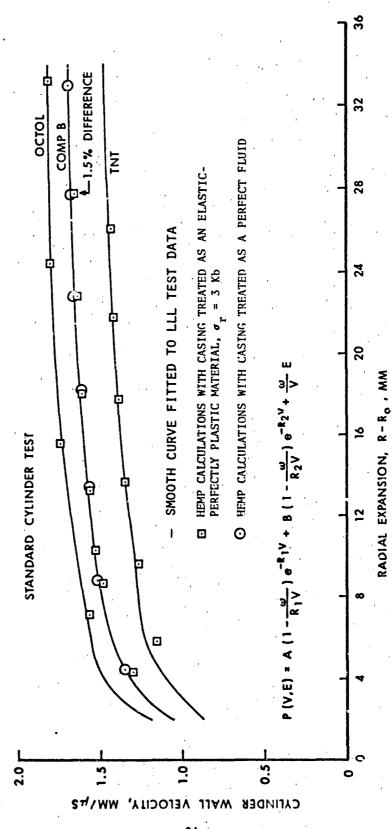
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Figure 7. Cylinder Expansion and Measured Values of Expansion Ratios for Various Casing Materials at which Gas Leakage Occurs (Ref. 4)



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Figure 8. HEMP Calculations Using the JWL Equation of State of the Explosive as Indicated Above Compared with the Data from the LLL Standard Cylinder Test

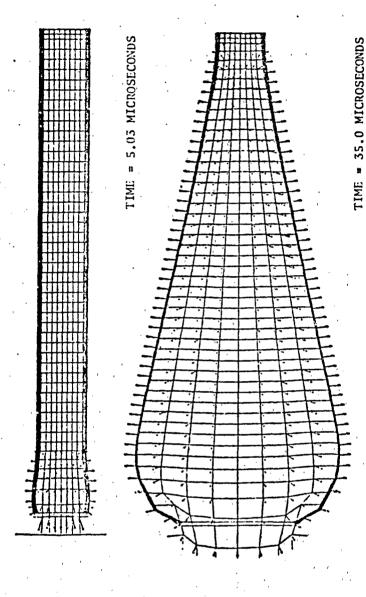
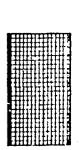
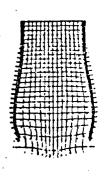


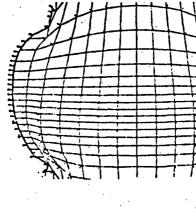
Figure 9. Finite-Difference Mesh Used in the HEMP Calculations of the LLL Standard Cylinder Test

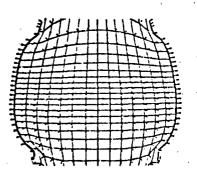


0 MICROSECONDS



25 MICROSECONDS

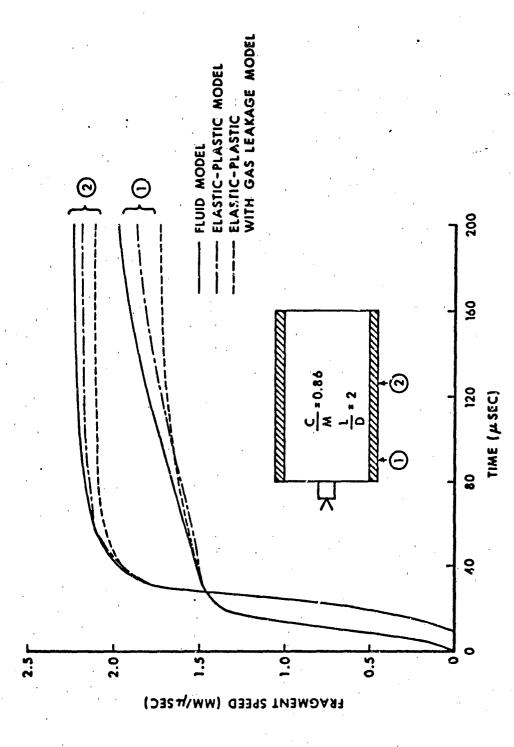




TIME * 75 MICROSECONDS

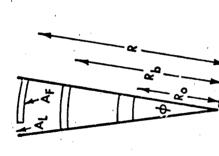
125 MICROSLCONDS

COMPUTER SIMULATION OF AN OPEN ENDED STEEL CYLINDER WITH AN OCTOL EXPLOSIVE TEST



Comparison of Fragment Speed Versus Time at Two Positions on the Casing Wall Using the Fluid, Elastic-Plastic and Elastic-Plastic with Gas Leakage Models

Figure 11.



Ro-INITIAL CASING RADIUS
Rb-RADIUS AT FRAGMENTATION
R - CURRENT RADIUS
AF= Rb \phi - FRAGMENT AREA/UNIT LENGTH
AL= (R-Rb) \ph-LEAKAGE AREA/UNIT LENGTH

FORCE ON FRAGMENT/UNIT LENGTH = $PA_F = P(\frac{A_F}{A_L + A_F}) (A_L + A_F)$ $= \frac{Rb}{R} P(A_L + A_F)$ = (MODIFIED PRESSURE) (AREA IN CODE)

Figure 12. Notation for Gas Leakage Model

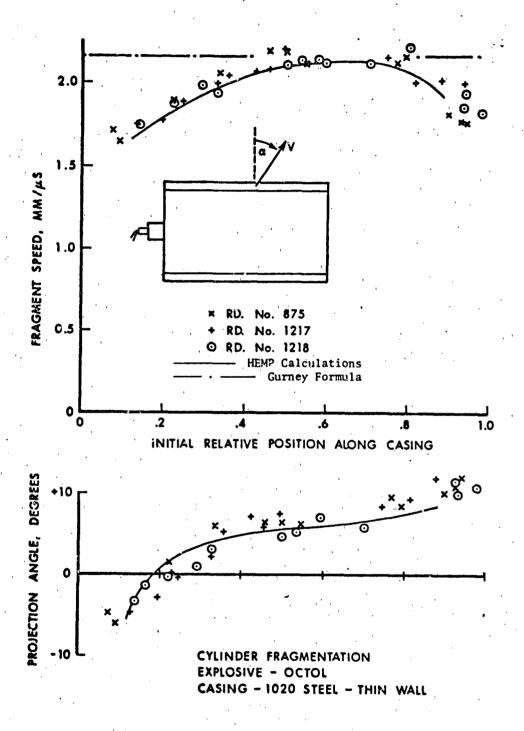


Figure 13. Comparison Between HEMP Calculations and Experimental Data of Fragment Speed and Projection Angle for Cylinders of L/D = 2, Octol, C/M = .8

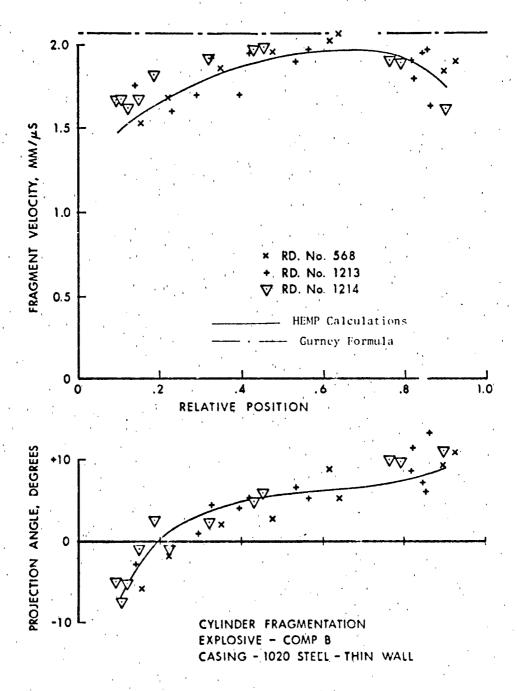


Figure 14. Comparison Between HEMP Calculations and Experimental Data of Fragment Speed and Projection Angle for Cylinders of L/D = 2, Composition B, G/M = .8

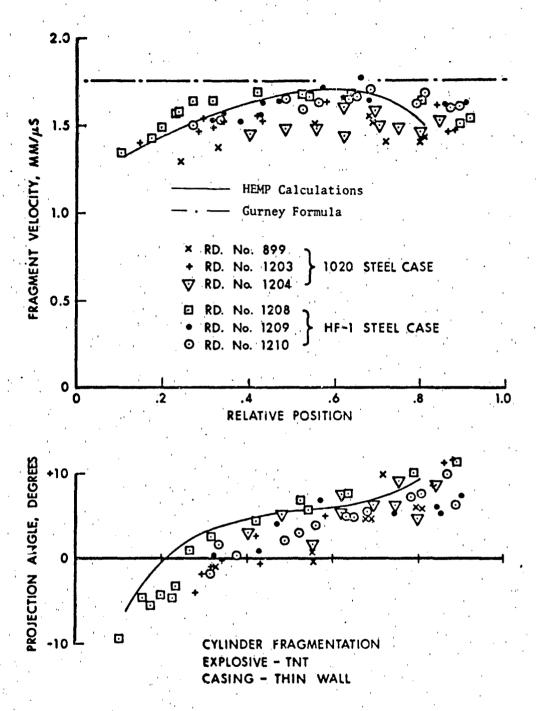


Figure 15. Comparison Between HEMP Calculations and Experimental Data of Fragment Speed and Projection Angle for Cylinders of L/D = 2, TNT, C/M = .8

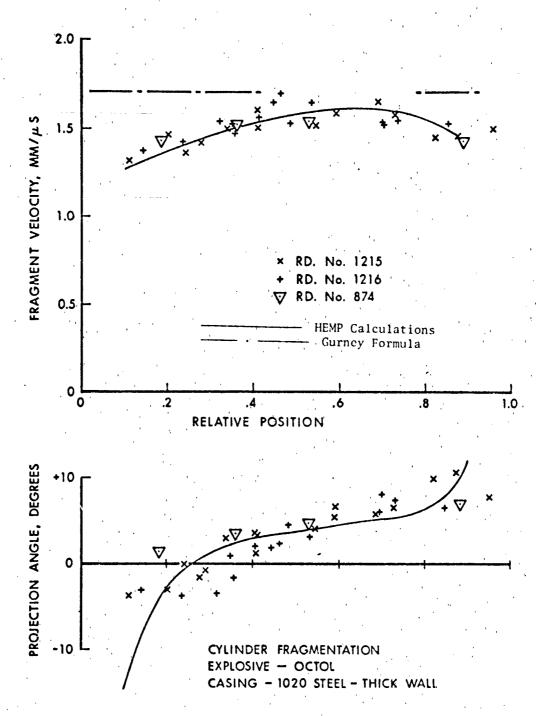


Figure 16. Comparison Between HEMP Calculations and Experimental Data of Fragment Speed and Projection Angle for Cylinders of L/D = 2, Octol, $C/M \approx .4$

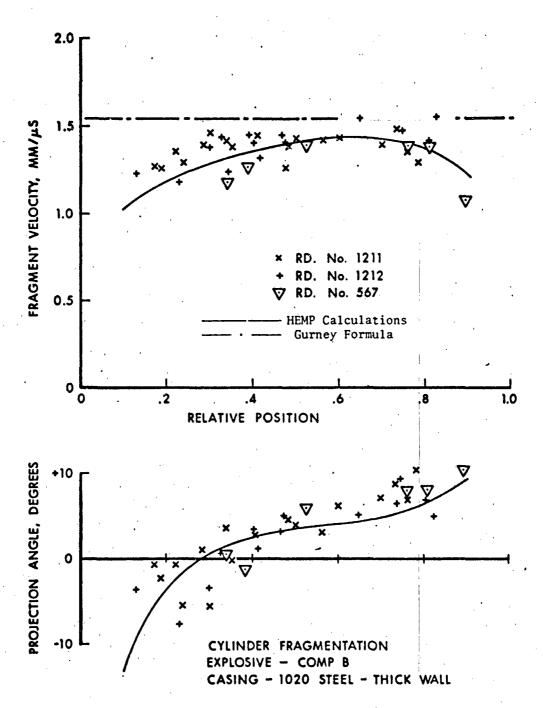
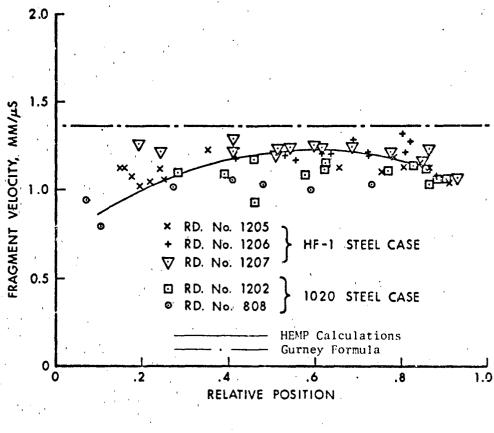


Figure 17. Comparison Between HEMP Calculations and Experimental Data of Fragment Speed and Projection Angle for Cylinders of L/D = 2, Composition B, C/M = .4



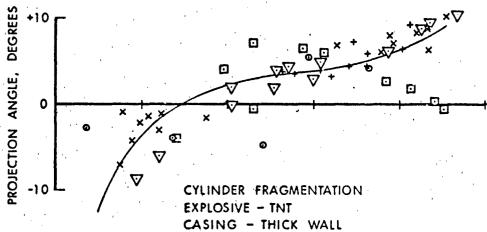


Figure 18. Comparison Between HEMP Calculations and Experimental Data of Fragment Speed and Projection Angle for Cylinders of L/D = 2, ΓNT , C/M = .4

105 MM, M1, HE PROJECTILE
FRAGMENT VELOCITY DISTRIBUTION
TEST DATA AND CALCULATED RESULTS

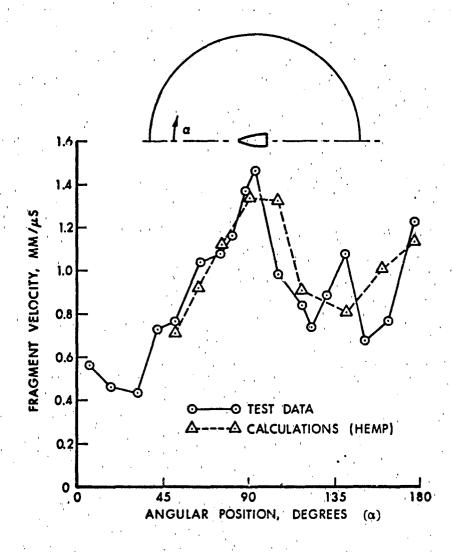


Figure 19. Comparison Between HEMP Calculations and Experimental
Data Obtained From Arena Tests - Fragment Speed Versus
Angular Position

105 MM, M1, HE PROJECTILE FRAGMENT MASS DISTRIBUTION TEST DATA AND CALCULATED RESULTS

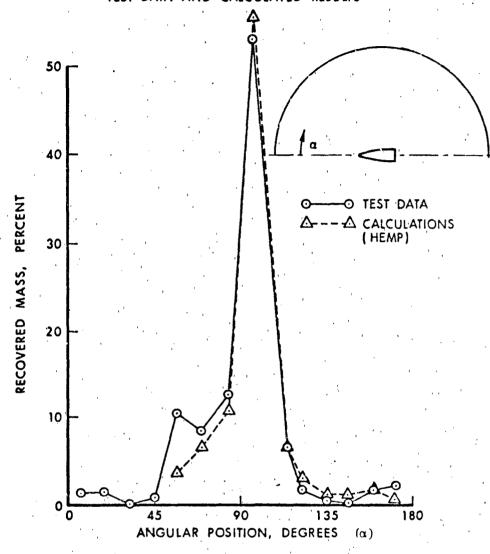


Figure 20. Comparison Between HEMP Calculations and Experimental Data Obtained From Arena Tests - Fragment Mass Versus Angular Position

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